Brave New World: Jet Production (and Disappearance) at the LHC Using the ATLAS Detector

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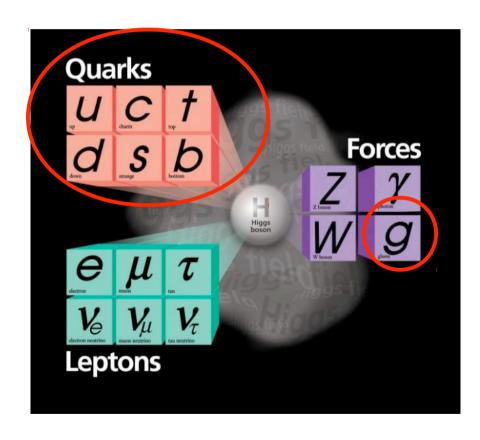
Outline

- Introduction to QCD and jet physics
- ATLAS Experiment at the LHC
 - Event displays of interesting jets & dijet events
 - Jet performance and calibration
- Inclusive jet and dijet cross-sections
 - Laboratory for perturbative QCD
- Searches for exotica using dijets
 - Resonances and contact interactions
- Jet quenching in lead ion collisions
- Conclusions and outlook

QCD and jet physics

Introduction

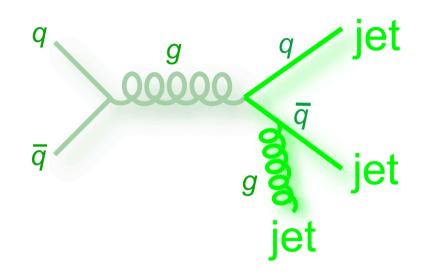
 Quantum chromodynamics (QCD) is the theory of the strong interaction, one of the three fundamental forces in the Standard Model

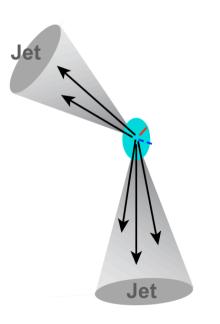


 $-\frac{1}{2}\partial_{\nu}g^{a}_{\mu}\partial_{\nu}g^{a}_{\mu} - g_{s}f^{abc}\partial_{\mu}g^{a}_{\nu}g^{b}_{\mu}g^{c}_{\nu} - \frac{1}{4}g^{2}_{s}f^{abc}f^{ade}g^{b}_{\mu}g^{c}_{\nu}g^{d}_{\mu}g^{e}_{\nu}$ $v_h^2 H^2 - \partial_\mu \phi^+ \partial_\mu \phi^- - M^2 \phi^+ \phi^- - \frac{1}{2} \partial_\mu \phi^0 \partial_\mu \phi^0 - \frac{1}{2e^2} M \phi^0 \phi^0 - \beta_0$ $W_{\nu}^{-}\partial_{\nu}W_{\mu}^{+})] - igs_{w}[\partial_{\nu}A_{\mu}(W_{\mu}^{+}W_{\nu}^{-} - W_{\nu}^{+}W_{\mu}^{-}) - A_{\nu}(W_{\mu}^{+}\partial_{\nu}W_{\mu}^{-})]$ $W_{\mu}^{-}\partial_{\nu}W_{\mu}^{+}) + A_{\mu}(W_{\nu}^{+}\partial_{\nu}W_{\mu}^{-} - W_{\nu}^{-}\partial_{\nu}W_{\mu}^{+})] - \frac{1}{2}g^{2}W_{\mu}^{+}W_{\mu}^{-}W_{\nu}^{+}W_{\nu}^{-} +$ $\frac{1}{2}g^{2}W_{\mu}^{+}W_{\nu}^{-}W_{\mu}^{+}W_{\nu}^{-} + g^{2}c_{w}^{2}(Z_{\mu}^{0}W_{\mu}^{+}Z_{\nu}^{0}W_{\nu}^{-} - Z_{\mu}^{0}Z_{\mu}^{0}W_{\nu}^{+}W_{\nu}^{-}) +$ $g^2 s_w^2 (A_\mu W_\mu^+ A_\nu W_\nu^- - A_\mu A_\mu W_\nu^+ W_\nu^-) + g^2 s_w c_w [A_\mu Z_\nu^0 (W_\mu^+ W_\nu^- W_{\nu}^{+}W_{\mu}^{-}$) $-2A_{\mu}Z_{\mu}^{0}W_{\nu}^{+}W_{\nu}^{-}$] $-g\alpha[H^{3}+H\phi^{0}\phi^{0}+2H\phi^{+}\phi^{-}]$ - $\frac{1}{8}g^{2}\alpha_{h}[H^{4}+(\phi^{0})^{4}+4(\phi^{+}\phi^{-})^{2}+4(\phi^{0})^{2}\phi^{+}\phi^{-}+4H^{2}\phi^{+}\phi^{-}+2(\phi^{0})^{2}H^{2}]$ $gMW_{\mu}^{+}W_{\mu}^{-}H - \frac{1}{2}g\frac{M}{c^{2}}Z_{\mu}^{0}Z_{\mu}^{0}H - \frac{1}{2}ig[W_{\mu}^{+}(\phi^{0}\partial_{\mu}\phi^{-} - \phi^{-}\partial_{\mu}\phi^{0}) W_{\mu}^{-}(\phi^{0}\partial_{\mu}\phi^{+} - \phi^{+}\partial_{\mu}\phi^{0})] + \frac{1}{2}g[W_{\mu}^{+}(H\partial_{\mu}\phi^{-} - \phi^{-}\partial_{\mu}H) - W_{\mu}^{-}(H\partial_{\mu}\phi^{+} - \phi^{-}\partial_{\mu}H)] + \frac{1}{2}g[W_{\mu}^{+}(H\partial_{\mu}\phi^{-} - \phi^{-}\partial_{\mu}H) - W_{\mu}^{-}(H\partial_{\mu}\phi^{+} - \phi^{-}\partial_{\mu}H)]$ $[\phi^{+}\partial_{\mu}H)] + \frac{1}{2}g\frac{1}{s_{-}}(Z_{\mu}^{0}(H\partial_{\mu}\phi^{0} - \phi^{0}\partial_{\mu}H) - ig\frac{s_{-}^{2}}{s_{-}}MZ_{\mu}^{0}(W_{\mu}^{+}\phi^{-} - W_{\mu}^{-}\phi^{+}) +$ $igs_w MA_\mu (W_\mu^+ \phi^- - W_\mu^- \phi^+) - ig \frac{1-2c_u^2}{2c_u} Z_\mu^0 (\phi^+ \partial_\mu \phi^- - \phi^- \partial_\mu \phi^+) +$ $igs_w A_\mu (\phi^+ \partial_\mu \phi^- - \phi^- \partial_\mu \phi^+) - \frac{1}{4} g^2 W_\mu^+ W_\mu^- [H^2 + (\phi^0)^2 + 2\phi^+ \phi^-] \frac{1}{4}g^2\frac{1}{c^2}Z_{\mu}^0Z_{\mu}^0[H^2 + (\phi^0)^2 + 2(2s_w^2 - 1)^2\phi^+\phi^-] - \frac{1}{2}g^2\frac{s_u^2}{c}Z_{\mu}^0\phi^0(W_{\mu}^+\phi^- +$ $W_{\mu}^{-}\phi^{+}$) $-\frac{1}{2}ig^{2}\frac{s_{w}^{2}}{c_{w}}Z_{\mu}^{0}H(W_{\mu}^{+}\phi^{-}-W_{\mu}^{-}\phi^{+})+\frac{1}{2}g^{2}s_{w}A_{\mu}\phi^{0}(W_{\mu}^{+}\phi^{-}+$ $W_{\mu}^{-}\phi^{+}) + \frac{1}{2}ig^{2}s_{w}A_{\mu}H(W_{\mu}^{+}\phi^{-} - W_{\mu}^{-}\phi^{+}) - g^{2}\frac{s_{w}}{c_{w}}(2c_{w}^{2} - 1)Z_{\mu}^{0}A_{\mu}\phi^{+}\phi^{-} - W_{\mu}^{-}\phi^{+})$ $g^1 s_w^2 A_\mu A_\mu \phi^+ \phi^- - \bar{e}^\lambda (\gamma \partial + m_e^\lambda) e^\lambda - \bar{\nu}^\lambda \gamma \partial \nu^{\lambda} - \bar{u}_i^\lambda (\gamma \partial + m_u^\lambda) u_i^\lambda \bar{d}_{i}^{\lambda}(\gamma \partial + m_{d}^{\lambda})d_{i}^{\lambda} + igs_{w}A_{\mu}[-(\bar{e}^{\lambda}\gamma^{\mu}e^{\lambda}) + \frac{2}{3}(\bar{u}_{i}^{\lambda}\gamma^{\mu}u_{i}^{\lambda}) - \frac{1}{3}(\bar{d}_{i}^{\lambda}\gamma^{\mu}d_{i}^{\lambda})] +$ $\frac{ig}{4c_{-}}Z_{\mu}^{0}[(\bar{\nu}^{\lambda}\gamma^{\mu}(1+\gamma^{5})\nu^{\lambda}) + (\bar{e}^{\lambda}\gamma^{\mu}(4s_{w}^{2}-1-\gamma^{5})e^{\lambda}) + (\bar{u}_{j}^{\lambda}\gamma^{\mu}(\frac{4}{3}s_{w}^{2}-1)e^{\lambda})]$ $(1 - \gamma^5)u_j^{\lambda}$ + $(\bar{d}_j^{\lambda}\gamma^{\mu}(1 - \frac{8}{3}s_w^2 - \gamma^5)d_j^{\lambda})$] + $\frac{ig}{2\sqrt{2}}W_{\mu}^+[(\bar{\nu}^{\lambda}\gamma^{\mu}(1 + \gamma^5)e^{\lambda})$ + $(\bar{u}_i^{\lambda}\gamma^{\mu}(1+\gamma^5)C_{\lambda\kappa}d_j^{\kappa})] + \frac{ig}{2\sqrt{2}}W_{\mu}^-[(\bar{e}^{\lambda}\gamma^{\mu}(1+\gamma^5)\nu^{\lambda}) + (\bar{d}_j^{\kappa}C_{\lambda\kappa}^{\dagger}\gamma^{\mu}(1+\gamma^5)\nu^{\lambda})]$ $\gamma^5 u_j^{\lambda}$] + $\frac{ig}{2\sqrt{2}} \frac{m_c^{\lambda}}{M} [-\phi^+(\bar{\nu}^{\lambda}(1-\gamma^5)e^{\lambda}) + \phi^-(\bar{e}^{\lambda}(1+\gamma^5)\nu^{\lambda})] \frac{g}{2}\frac{m_e^{\lambda}}{M}[H(\bar{e}^{\lambda}e^{\lambda}) + i\phi^0(\bar{e}^{\lambda}\gamma^5e^{\lambda})] + \frac{ig}{2M\sqrt{2}}\phi^+[-m_d^{\kappa}(\bar{u}_j^{\lambda}C_{\lambda\kappa}(1-\gamma^5)d_j^{\kappa}) +$ $m_u^{\lambda}(\bar{u}_j^{\lambda}C_{\lambda\kappa}(1+\gamma^5)d_j^{\kappa}] + \frac{ig}{2M\sqrt{2}}\phi^-[m_d^{\lambda}(\bar{d}_j^{\lambda}C_{\lambda\kappa}^{\dagger}(1+\gamma^5)u_j^{\kappa}) - m_u^{\kappa}(\bar{d}_j^{\lambda}C_{\lambda\kappa}^{\dagger}(1-\gamma^5)u_j^{\kappa})]$ $\gamma^5 u_j^{\kappa} = \frac{g}{2} \frac{m_{\gamma}^{\lambda}}{M} H(\bar{u}_j^{\lambda} u_j^{\lambda}) - \frac{g}{2} \frac{m_{\dot{q}}^{\lambda}}{M} H(\bar{d}_j^{\lambda} d_j^{\lambda}) + \frac{ig}{2} \frac{m_{\dot{q}}^{\lambda}}{M} \phi^0(\bar{u}_j^{\lambda} \gamma^5 u_j^{\lambda}) - \frac{g}{2} \frac{m_{\dot{q}}}{M} \phi^0(\bar{u}$ $\frac{ig}{2} \frac{m_{\lambda}^{\lambda}}{M} \phi^{0}(\bar{d}_{i}^{\lambda} \gamma^{5} d_{i}^{\lambda}) + \bar{X}^{+}(\partial^{2} - M^{2})X^{+} + \bar{X}^{-}(\partial^{2} - M^{2})X^{-} + \bar{X}^{0}(\partial^{2} - M^{2})X^{-}$ $\frac{M^2}{c^2}$ $X^0 + \bar{Y}\partial^2 Y + igc_wW^+_\mu(\partial_\mu \bar{X}^0 X^- - \partial_\mu \bar{X}^+ X^0) + igs_wW^+_\mu(\partial_\mu \bar{Y} X^- - \partial_\mu \bar{X}^+ X^0)$ $\partial_{\mu}\bar{X}^{+}Y$) + $igc_{w}W_{\mu}^{-}(\partial_{\mu}\bar{X}^{-}X^{0} - \partial_{\mu}\bar{X}^{0}X^{+}) + igs_{w}W_{\mu}^{-}(\partial_{\mu}\bar{X}^{-}Y - \partial_{\mu}\bar{X}^{0}X^{+})$ $\partial_{\mu} \bar{Y} X^{+}$) + $igc_{w}Z^{0}_{\mu}(\partial_{\mu} \bar{X}^{+} X^{+} - \partial_{\mu} \bar{X}^{-} X^{-}) + igs_{w}A_{\mu}(\partial_{\mu} \bar{X}^{+} X^{+} - \partial_{\mu} \bar{X}^{-} X^{-})$ $\partial_{\mu}\bar{X}^{-}X^{-}) - \frac{1}{2}gM[\bar{X}^{+}X^{+}H + \bar{X}^{-}X^{-}H + \frac{1}{c!}\bar{X}^{0}X^{0}H] +$ $\frac{1-2c_w^2}{2c_w}igM[\bar{X}^+X^0\phi^+ - \bar{X}^-X^0\phi^-] + \frac{1}{2c_w}igM[\bar{X}^0X^-\phi^+ - \bar{X}^0X^+\phi^-] +$ $igMs_w[\bar{X}^0X^-\phi^+ - \bar{X}^0X^+\phi^-] + \frac{1}{2}igM[\bar{X}^+X^+\phi^0 - \bar{X}^-X^-\phi^0]$

QCD & Jets at the LHC

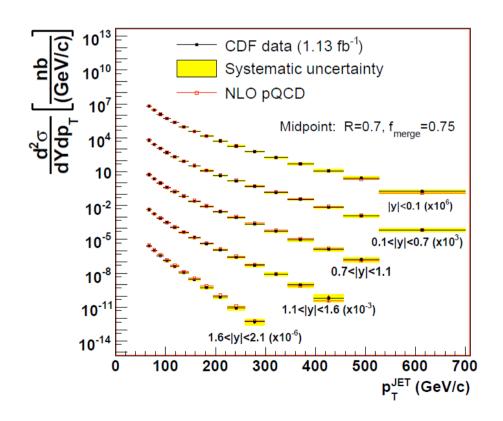
- QCD is ideal candidate for early LHC physics
- Strongly coupled theory → large production cross-sections for jets (collimated flows of hadrons)
- Relatively "simple" final state topologies with close connection to detector performance, e.g. jet calibration
- Tests of Standard Model
 - Probes of NLO perturbative QCD
 - Also sensitive to non-perturbative effects
- Jets are perfect to discover new physics
 - Large production crosssections and small backgrounds
 - Highest sensitivity to new physics with early data
- Background to many new physics channels

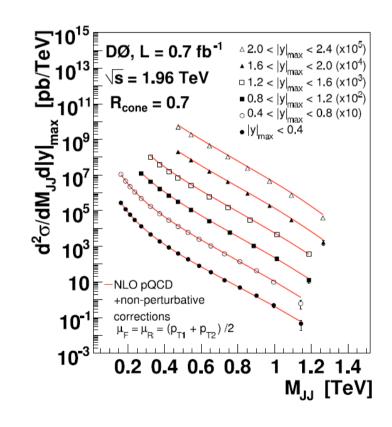




Existing jet measurements

- Many jet measurements using Tevatron ppbar collider at Fermilab
- Inclusive jet pT spectrum extends to p_⊤ of 700 GeV (left: CDF)
- Dijet mass spectrum up to 1.4 TeV (right: D0)

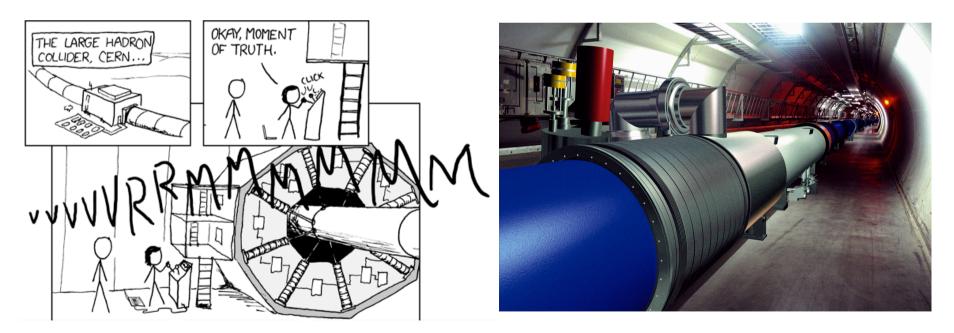




Enter the LHC...

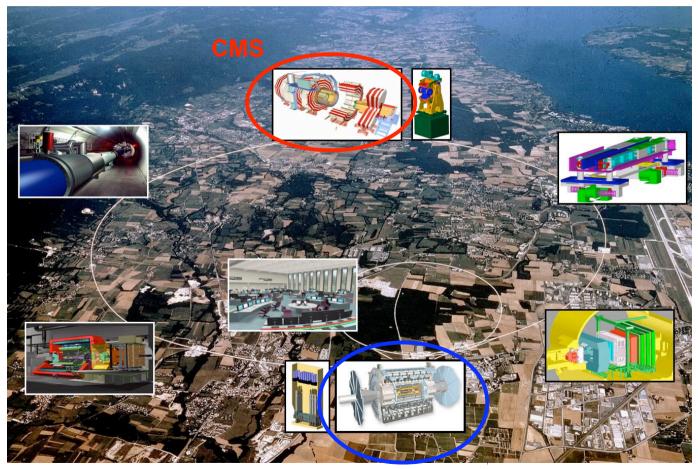
High-p_T jet factory

- CERN Large Hadron Collider (LHC) in Geneva, Switzerland has a monopoly on high-p_⊤ jet production
- Highest p_⊤ jets and largest dijet masses ever produced
 - pp collisions at sqrt(s) = 7 TeV (design of 14 TeV)
 - 3.5 times larger center-of-mass energy than Tevatron



CERN Large Hadron Collider

- ATLAS and CMS are two general-purpose detectors at the LHC
- Relatively small statistics, but accelerating quickly



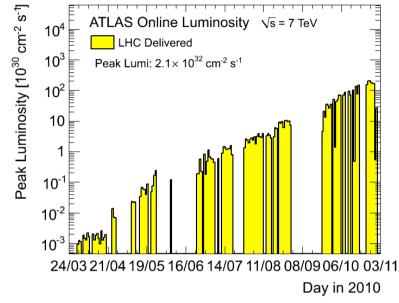
LHC proton beams at 3.5 TeV

- After initial hiccup during first beam in November 2008, performance of the LHC has been outstanding
 - First collisions at 900 GeV in late 2009
 - Collisions at 7 TeV in early 2010
 - Intense media coverage led to unintended comedy at times

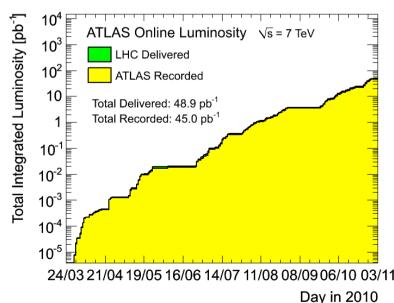


Luminosity of the machine

- Instantaneous luminosity increased by 5 orders of magnitude since March
 - Peak instantaneous luminosity in $2010 = 2.1 \times 10^{32} \text{ cm}^2 \text{ s}^4$



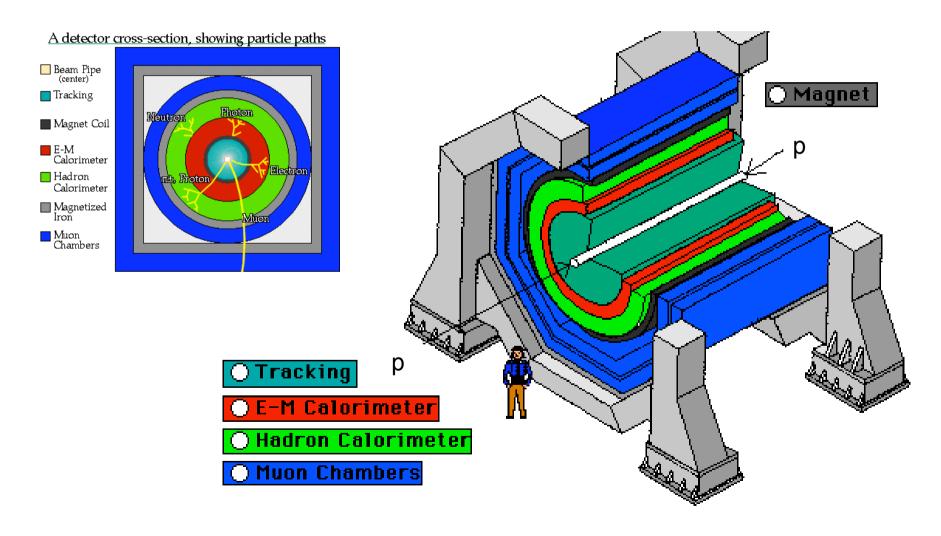
- Dataset
 - ICHEP (July): 17 nb⁻¹
 - These results: 3 pb⁻¹
 - Full 2010 dataset (Winter): 45 pb⁻¹



The ATLAS experiment

ATLAS Detector (schematic)

Schematic of a general-purpose detector



Muon Spectrometer ($|\eta|$ <2.7) : air-core toroids with gas-based muon chambers Muon trigger and measurement with momentum resolution < 10% up to E_u ~ 1 TeV

ATLAS Detector

Muon Detectors Tile Calorimeter Liquid Argon Calorimeter

3-level trigger reducing the rate from 40 MHz to ~200 Hz

Length: ~ 46 m

Radius : ~ 12 m

Weight: ~ 7000 tons ~108 electronic channels

3000 km of cables

Inner Detector ($|\eta|$ <2.5, B=2T): Si Pixels, Si strips, Transition Radiation detector (straws) Precise tracking and vertexing, e/ π separation Momentum resolution:

 $\sigma/p_{T} \sim 3.8 \times 10^{-4} p_{T} (GeV) \oplus 0.015$

Toroid Magnets Solenoid Magnet SCT Tracker Pixel Detector TRT Tracker

EM calorimeter: Pb-LAr Accordion

e/γ trigger, identification and measurement

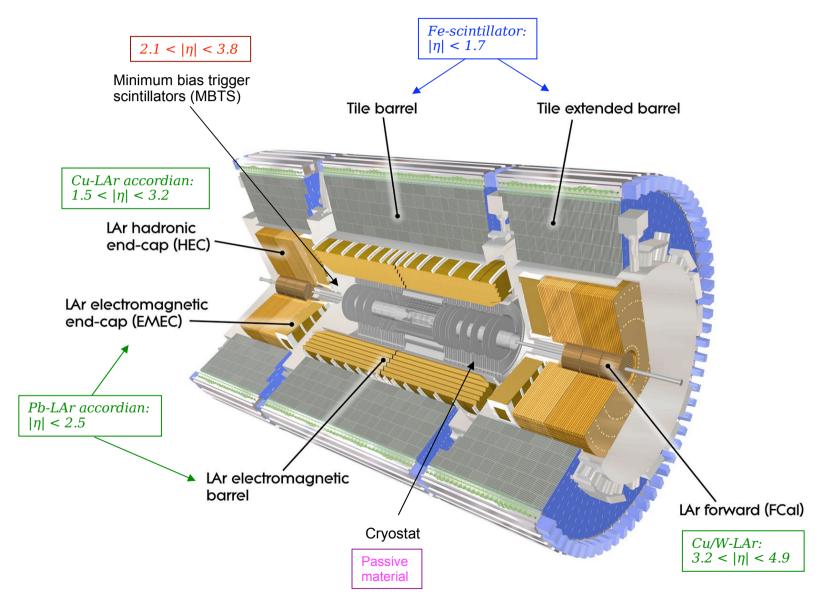
E-resolution: $\sigma/E \sim 10\%/\sqrt{E}$

HAD calorimetry ($|\eta|$ <5): segmentation, hermeticity Fe/scintillator Tiles (central), Cu/W-LAr (fwd) Trigger and measurement of jets and missing E_T

E-resolution: $\sigma/E \sim 50\%/\sqrt{E} \oplus 0.03$

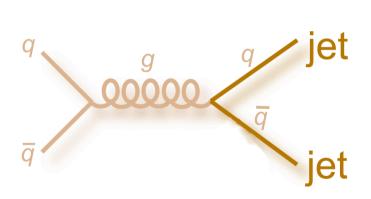
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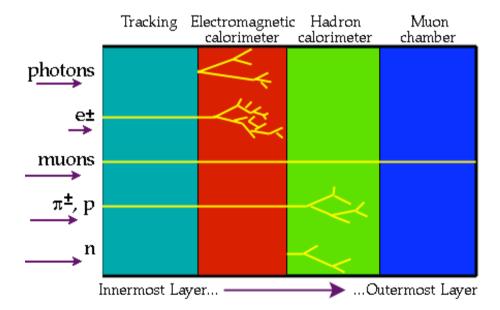
ATLAS Calorimeter



Inclusive jet production

- Leading order diagram in SM is dijet production
 - Deposits energy primarily in hadronic calorimeter
- Define jets inclusively as: "calorimeter energy deposition"
 - Backgrounds from electrons, photons, etc are negligible because QCD cross-sections are so large





"QCD in Year One": Then and Now

- At the ATLAS Workshop of the Americas at NYU in 2009, I described in the QCD plenary the various jet measurements we expected to make with the first year of data.
- After ~8 months of *pp* collisions at 7 TeV, where do we stand now?
 - Inclusive jet p_⊤
 - Dijet mass spectrum
 - Dijet χ angular distribution
 - Dijet Δφ decorrelation
 - Dijets with rapidity gaps
 - Multi-jets■
 - W/Z + jets ⊠
 - Jet shapes
 - Searches for exotica
 - Plus a bonus (jet quenching) in lead ion collisions!

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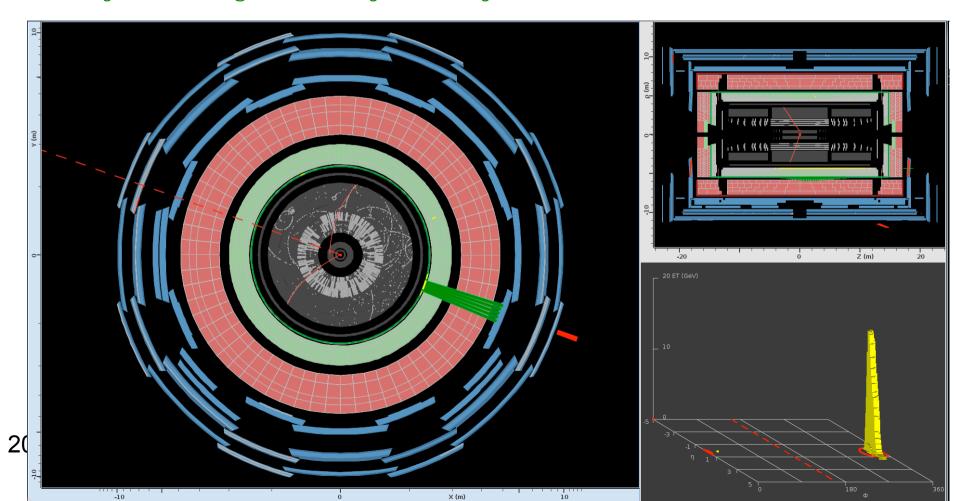
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Will focus on these analyses today

What do jets look like in the ATLAS calorimeter?

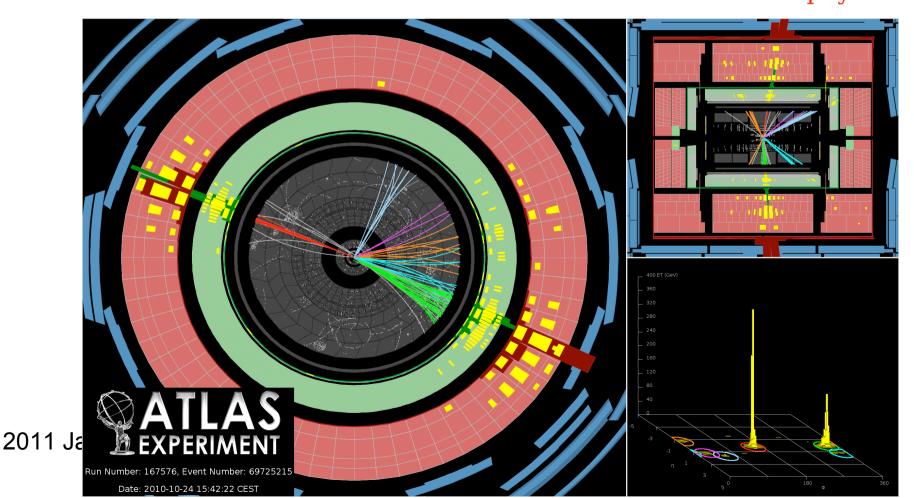
What a high p_T jet does NOT look like

- Fake jet with p_{T} of 1.1 TeV
 - Beam halo particle showers longitudinally in EM presampler so its energy receives large weight
- Missing ET = 1.3 TeV (exactly opposite in φ to leading jet)
- Add jet cleaning cuts to reject fake jets like this



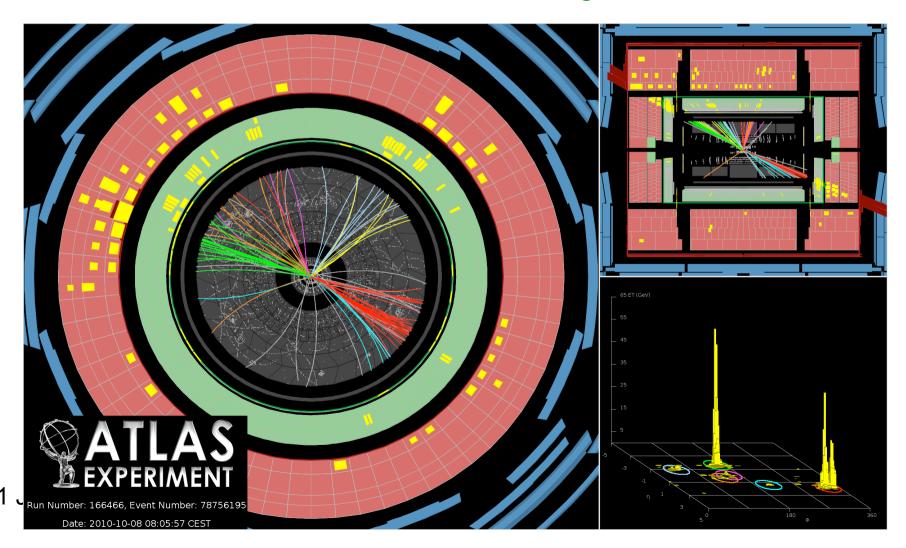
High p_T jet (1.3 TeV) and large central dijet mass (2.6 TeV)

- Two central, well-measured jets in |y| < 0.8 with very high- p_T
 - 1st jet: pT = 1.3 TeV, $\eta = 0.2$, $\phi = 2.8$
 - 2nd jet: pT = 1.2 TeV, $\eta = 0.0$, $\varphi = -0.5$
- These events are the most interesting because they have large momentum transfer and s-channel is more sensitive to new physics



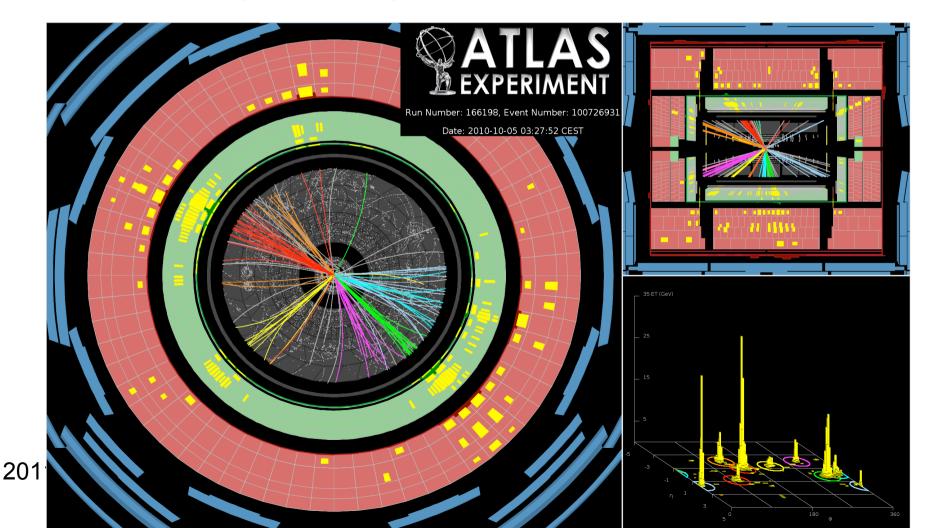
Large dijet mass (3.7 TeV)

- Dijet mass generated by rapidity separation of forward jets
 - 1st jet: pT = 670 GeV, $\eta = 1.9$, $\varphi = -0.5$
 - 2nd jet: pT = 610 GeV, η = -1.6, φ = 2.8
- Characteristic of t-channel and u-channel scattering in QCD



High jet multiplicity event (8 jets above pT = 60 GeV)

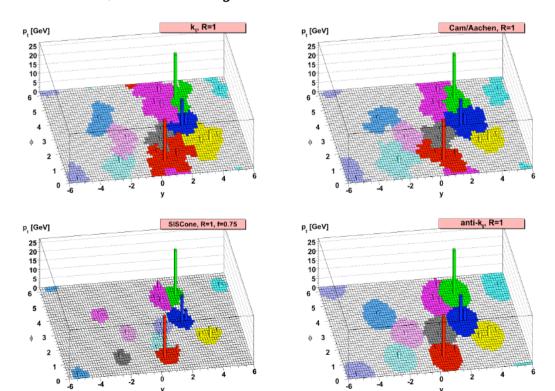
- Eight jets between 60-300 GeV (associated to single primary vertex)
 - 1st jet: pT = 290 GeV, η = -0.9, φ = 2.7
 - 2nd jet: pT = 220 GeV, $\eta = 0.3$, $\varphi = -0.7$
- Exemplifies large amount of gluon radiation



Jet trigger, reconstruction, selection, and calibration

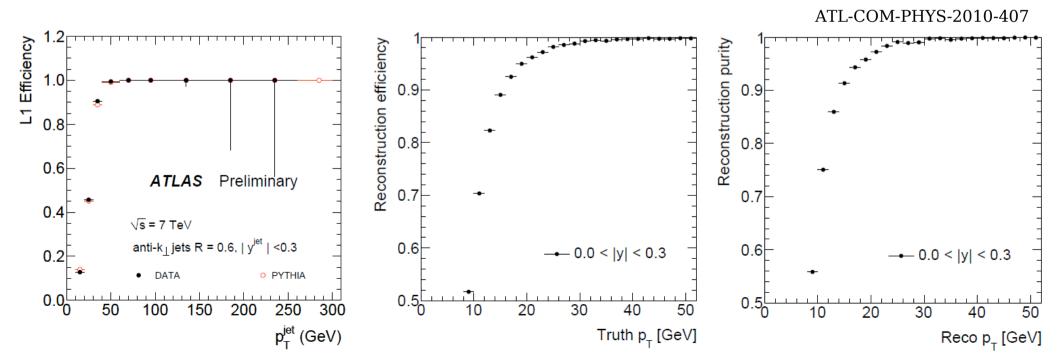
Jet reconstruction

- Jets reconstructed offline from calorimeter energy clusters using anti- $k_{_{\!\!T}}$ jet algorithm with clustering parameter R=0.6
 - Topological clusters are formed from seeds of calorimeter cells above noise thresholds to reject noise
 - Iteratively add neighbors farther away
 - Anti-k, is an infrared and collinear safe jet algorithm
 - Stable against underlying event and pileup
 - Uniform, cone-like jets



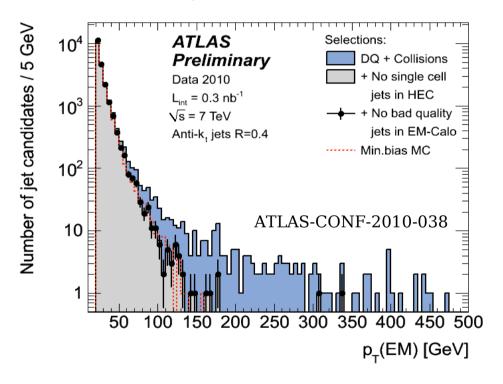
Jet trigger & reconstruction

- Require jet trigger fired
 - Trigger uses simpler and faster algorithms than offline reconstruction
- Lowest threshold jet trigger efficiency (left) is above 99% for $p_{\scriptscriptstyle T} > 60 \; \text{GeV}$
- Jet reconstruction efficiency (center) & purity (right) are each above 99% for $p_{\scriptscriptstyle T} > 30~\text{GeV}$
- \rightarrow Restrict leading (sub-leading) jet to $p_T > 60 \text{ GeV}$ (30 GeV)



Jet cleaning

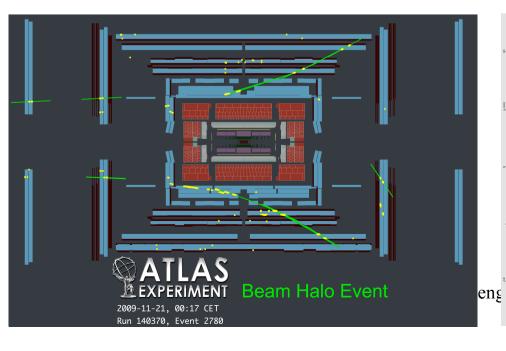
- "Bad" jet definition is based on cuts designed to remove:

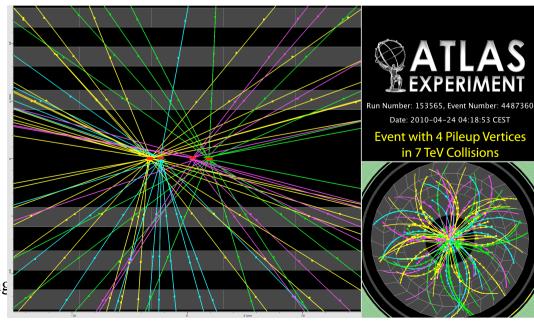


- Noisy cells in the hadronic endcap calorimeter
- Coherent noise in the electromagnetic calorimeter
- Large out-of-time energy depositions, e.g. from cosmic ray muons

Vertex requirement

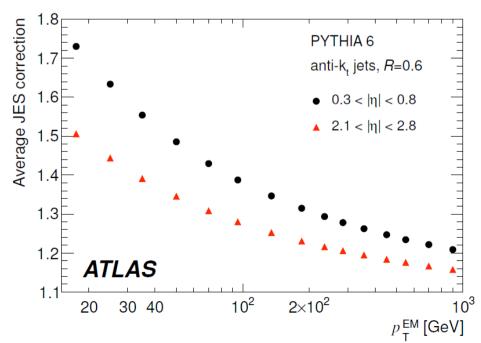
- Require at least one vertex reconstructed with at least 5 tracks and |z|<10cm (luminous region is ~5cm in early data, somewhat wider later)
 - Suppresses contamination from non-collision backgrounds such as beam halo and beam gas
 - 99% efficient for early data
- Effect of pileup is accounted for in absolute JES uncertainty
- |z| vertez cut also helps to ensure that jets are well-measured because they originate from near the detector center

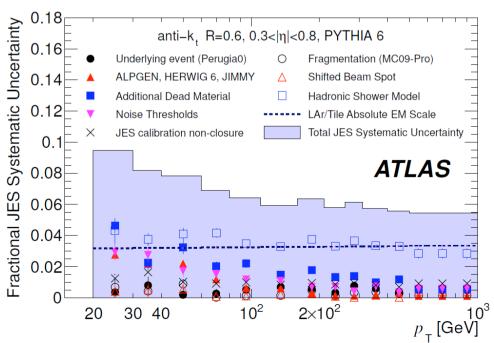




Jet energy scale (JES) calibration

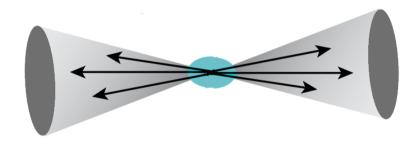
- - MC-based calibration as a function of jet $p_{\scriptscriptstyle T}$ and η
 - Non-linearity for hadrons due to non-compensation from nuclear interactions as well as dead material upstream
- Absolute JES uncertainty derived via systematic variation of parameters in MC
 - Conservative uncertainty within 7% for central jets with $p_T > 60 \text{ GeV}$ (dominant systematic uncertainty for most jet measurements)
 - Includes pileup uncertainty estimated using tower energy density





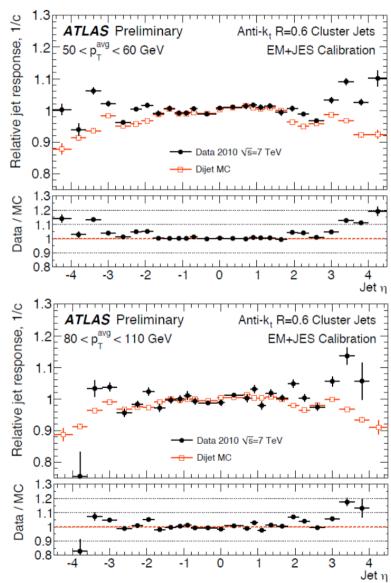
Relative jet energy scale

- Absolute JES assumed fully correlated in pseudorapidity
- Unlike single jets, dijets can span a large pseudorapidity
 → dependence also on the relative JES



- Dijet balance indicates relative JES of 5% within |y|<2.8
- Diverges to 15% in 2.8 < |y| < 4.5
- Thus restrict to |y|<2.8 when possible

ATLAS-CONF-2010-055



Jet energy resolution

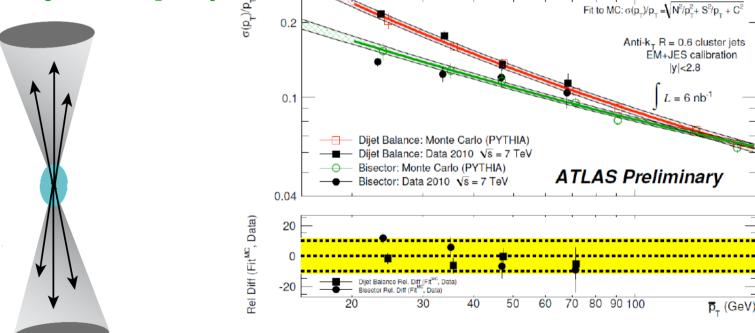
- Jet energy resolution in MC validated to ~14% via dijet balance and bisector methods (in-situ)
 - Control for systematic uncertainties from underlying event, initial state radiation, etc
 - Initial studies with more data show agreement within 10%

Jet angular resolution also verified at the level of 20% at very low p_T
 ATLAS-CONF-2010-054

Eric Feng, U. of Chicago

using track jets as a proxy

2011 Jan 4



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What the internal structure of jets looks like

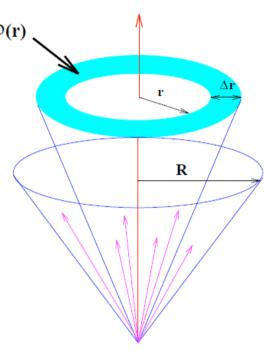
Differential jet shape: Method

• Differential jet shape in inclusive jet production:

$$\rho(r) = \frac{1}{\Delta r} \frac{1}{N^{jet}} \sum_{\text{jets}} \frac{p_T(r - \Delta r/2, r + \Delta r/2)}{p_T(0, R)}, \ 0 \le r \le R$$

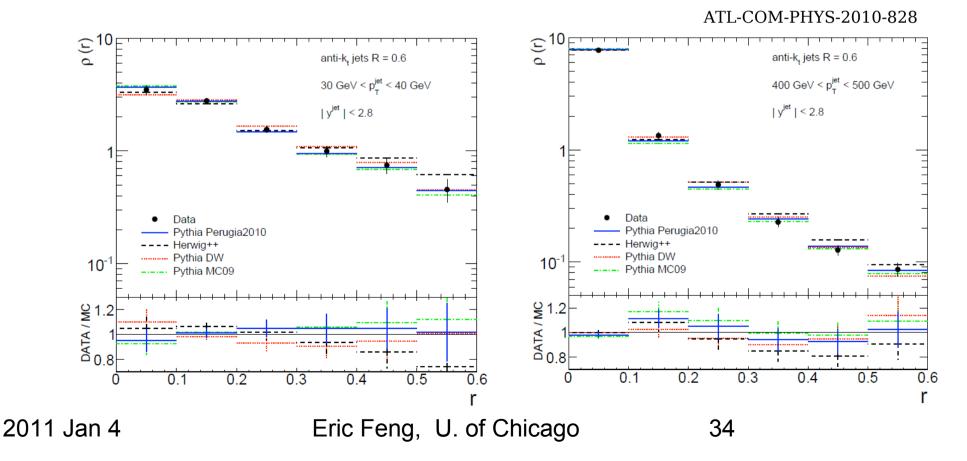


- Proportional to transverse momentum density inside the jet
- Relatively insensitive to the jet energy scale because jet p_⊤ normalized away
- Here p_T is computed as the scalar sum of transverse momenta of calorimeter energy clusters that lie within the annulus



Differential jet shapes: Results

- Jet shapes from calorimeter clusters used as basic validation
- Density of transverse momentum peaks at low r with most of jet p_T within r<0.3, indicating collimated flows of particles around jet axis
- Shifts to lower r for higher $p_{\scriptscriptstyle T}$ jets, indicating they are more collimated



How often are jets produced, and does this agree with NLO pQCD?

Unfolding

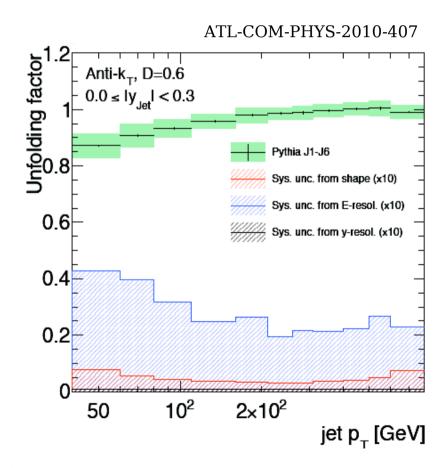
- Bin-by-bin unfolding back to hadron level, where all particles with lifetime > 10 ps (including muons and neutrinos) are included
- Correct for all detector effects: efficiencies, scales, and resolutions
 - Jet trigger efficiency
 - Vertex reconstruction efficiency
 - Jet reconstruction efficiency
 - Absolute and relative jet energy scale
 - Jet energy resolution
 - Jet angular resolution
- Strategy: Restrict kinematic region and bin the observable sufficiently coarsely so that corrections are both small and stable
- Do not correct for acceptance from $p_{\scriptscriptstyle T}$ and rapidity cuts
 - These are part of definition of the observable and no measurement is made outside this range

Unfolding systematics

Use Pythia MC09 to derive correction factors:

$$C = \sigma_{truth} / \sigma_{reco}$$

- Systematic uncertainty on this correction factor studied by:
 - Scaling up/down by JES uncertainty (dominant effect, not shown)
 - Worsening jet energy resolution by 15%
 - Smear angular resolution by 5%
 - Reweight cross-section to alter shape of distribution
- Also perform closure tests against Pythia 8, Herwig++, and Alpgen

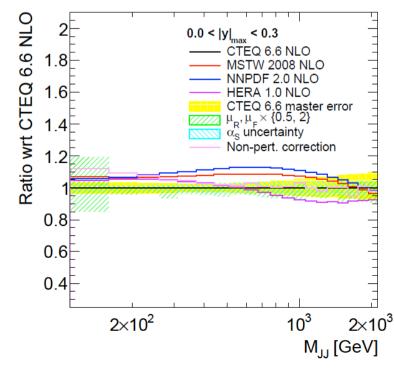


Theory prediction & uncertainties

- Monte Carlo used for observables with high jet multiplicity
 - LO 2->2: Pythia, Herwig
 - LO 2->N: Alpgen (matrix element generator)
- NLOJet++ used to calculate NLO pQCD prediction for low jet multiplicity
 - APPLgrid program used for efficient evaluation of uncertainties
 - **CTEQ 6.6 NLO PDF** is baseline, but also compared to MSTW 2008 NLO, NNPDF 2.0 NLO, HERA 1.0 NLO ATL-COM-PHYS-2010-407
- Renormalization scale μ_{R} varied by factor of 2 to account for neglected higher order terms
- Factorization scale μ_F varied by factor of 2 to account for scale separating matrix element from PDF evolution
- Strong coupling constant $\alpha_s(M_z)=0.118$ varied by 0.002 from world's best estimate

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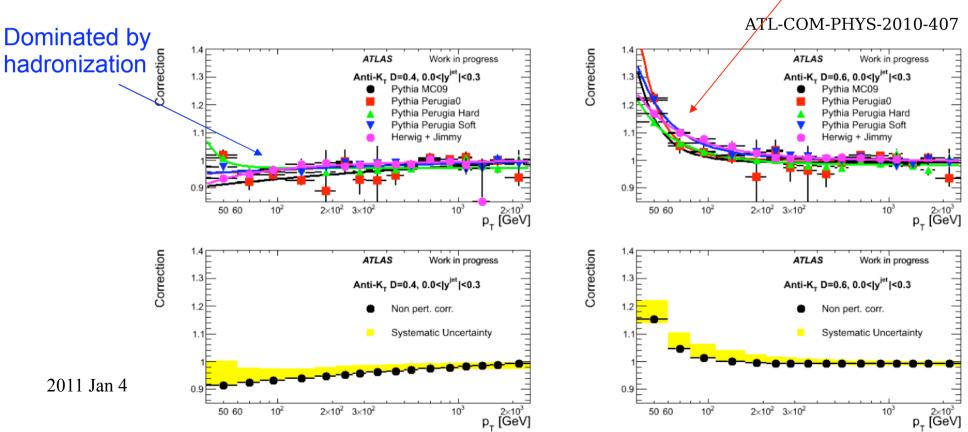


Non-perturbative correction

- Parton-level NLO calculation corrected for non-perturbative effects calculated using Rivet framework:
 - Hadronization and underlying event
 - Different dependence for R=0.4 (left) and 0.6 (right) → tune MC!

Systematic variations assessed using:

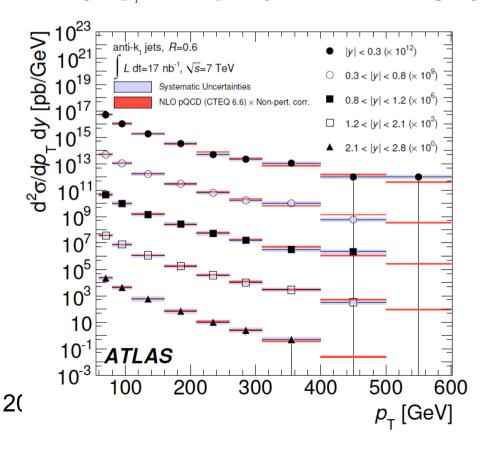
 Perugia0, Perugia hard, Perugia soft, Herwig

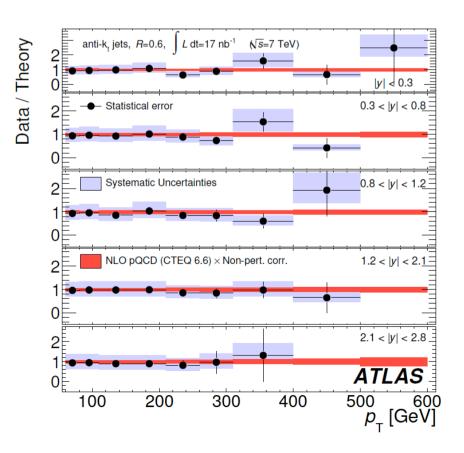


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Inclusive jet p_T cross-section

- ATLAS Collaboration. "Measurement of inclusive jet and dijet cross sections in protonproton collisions at 7 TeV centre-of-mass energy with the ATLAS detector." arXiv:1009.5908 [hep-ex]. Accepted by Eur. Phys. J. C.
- First cross-section measurements at a center-of-mass energy of 7 TeV of inclusive jet and dijet production, using 17 nb⁻¹ integrated luminosity
- Inclusive jet cross-section = Probability to observe a jet in a pp collision as a function of the jet p_{T} and rapidity (canonical single-jet observable)

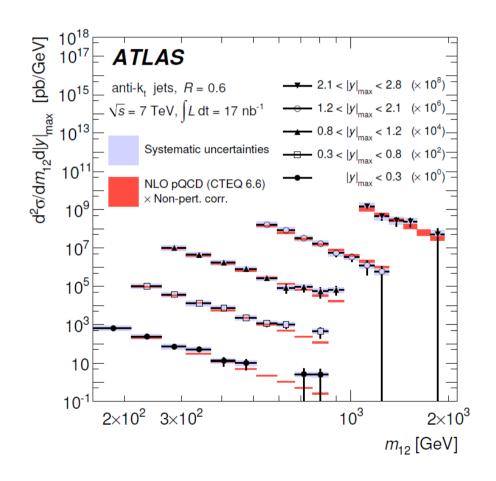




Dijet mass spectrum

Dijet mass = invariant mass of two leading jets measured in bins of maximum rapidity of two leading jets: $|y_{max}| = max(|y_1|, |y_2|)$

- Exotic resonances would tend to decay to high p_T central jets
- No discrepancy observed between data and NLO pQCD prediction

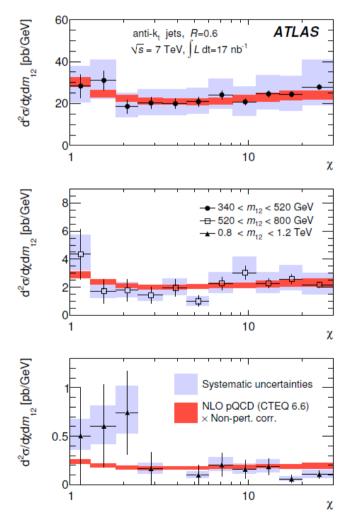


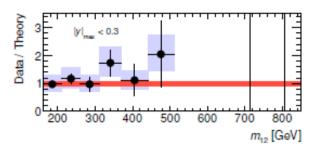
Dijet χ angular distribution

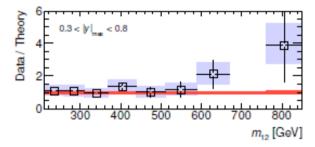
Dijet angular distribution: $\chi = \exp(|y_1 - y_2|) = (1 + \cos(\theta^*)) / (1 - \cos(\theta^*))$

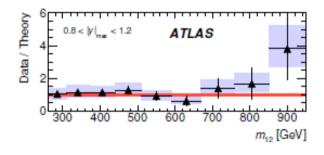
measured in bins of dijet mass m_{1,2}

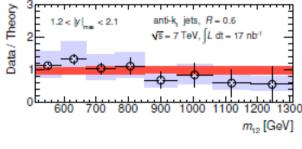
- Measures scattering angle θ^* in dijet center-of-mass frame
- Contact interactions such as compositeness would produce peak at low χ (small scattering angle), whereas QCD is ~ flat as we see
- Relatively insensitive to parton distribution functions because predictions for gg, qg, and qq subprocesses are similar
- No statistically significant deviations seen wrt pQCD prediction

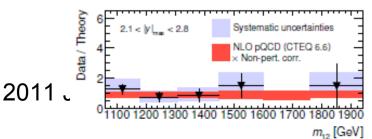








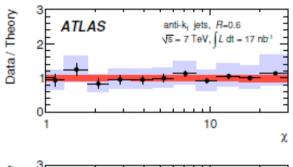


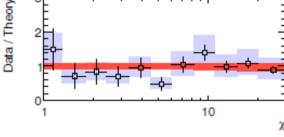


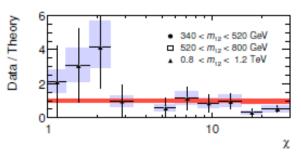
Dijet mass and χ: Data/theory ratio

NLO pQCD describes data well in all regions of phase space

- Both central and forward dijet mass follow the data (left)
- Angular distribution is as expected from pQCD at large dijet mass (right)



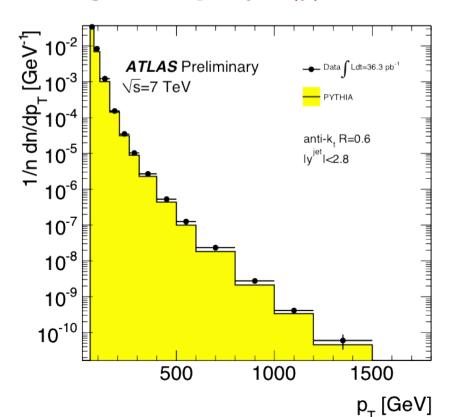


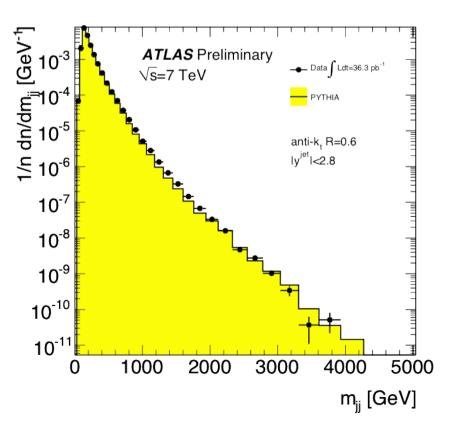


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Updated analysis with full 2010 dataset

- Full cross-section measurement with $\sim\!45~pb^{\cdot\!1}$ integrated lumi taken in 2010 is in progress
- Jets with transverse momentum up to 1.3 TeV are observed
- Invariant mass of two leading jets is observed up to 3.7 TeV
 - LO parton shower Monte Carlo provided to guide the eye



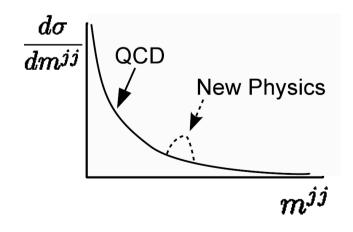


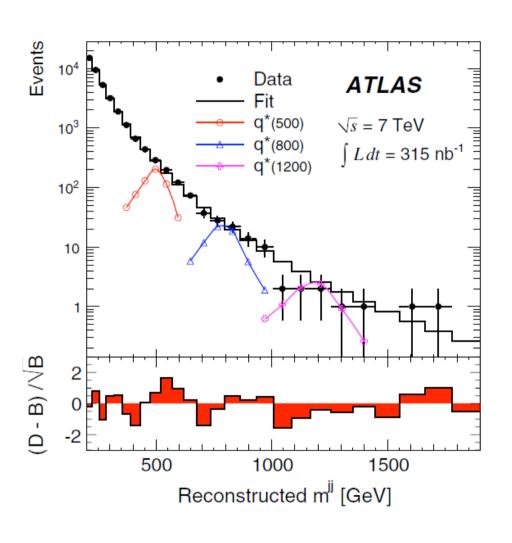
Are there any signs of new exotic physics at these short length scales?

Search for dijet resonance using the dijet mass spectrum

- There are direct extensions to searches for resonances decaying to two jets in final state (excited quarks, technirho, KK graviton, W'/Z', etc)
- Performed a search for bumps in the dijet mass spectrum by comparing it against a smooth continuum as predicted by QCD:

$$f(x) = p_0 \frac{(1-x)^{p_1}}{x^{p_2+p_3 \ln x}}, \quad x \equiv \frac{m^{jj}}{\sqrt{s}}$$

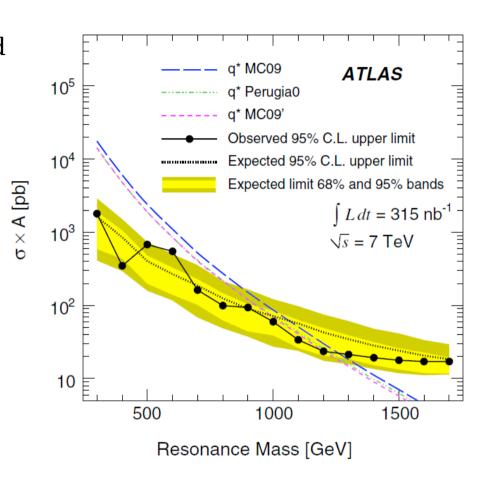




Limit on excited quarks

- No evidence of a resonance observed

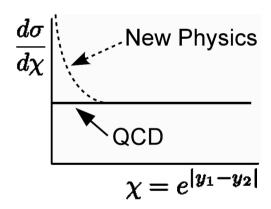
 → set new limit on dijet
 resonances arising from excited
 quarks
 - Bayesian approach used to set limit, cross-checked with frequentist method
- Excited quarks with mass up to 1.26 TeV excluded at 95% CL
 - Tevatron limit = 870 GeV
 - Systematics including JES (dominant), luminosity, background fit, etc also accounted for

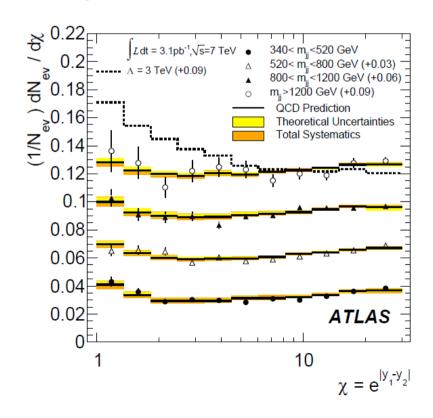


• ATLAS Collaboration. "Search for New Particles in Two-Jet Final States in 7 TeV Proton-Proton Collisions with the ATLAS Detector at the LHC." Phys Rev. Lett. 105, 161801 (2010)

Search for contact interactions with the dijet χ angular distribution

- Examined dijet $\chi = \exp(|y_1-y_2|)$ angular distribution for excess of low-angle scattering compared to NLO QCD prediction:
 - Normalized distribution mostly cancels out absolute JES uncertainty, leaving relative JES
- Any tail could be indication of contact interaction that might arise from quark compositeness, gravitational scattering, etc.

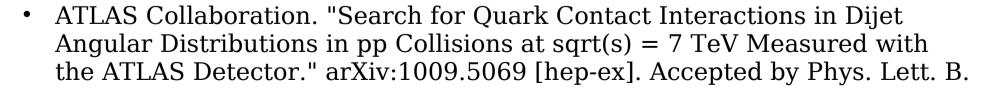


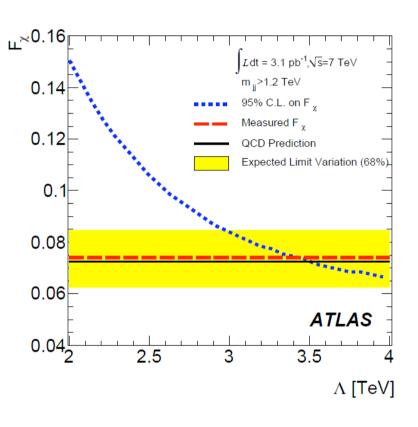


Limits on contact interactions

- No significant excess observed
- Set new limit on contact interactions that may arise from quark compositeness
 - Define F as fraction of events at low χ
 - Limit set using frequentistic method, cross-checked by Bayesian approach
- Using integrated luminosity of 3.1 pb $^{-1}$, compositeness scale $\Lambda < 3.4$ TeV excluded at 95% CL



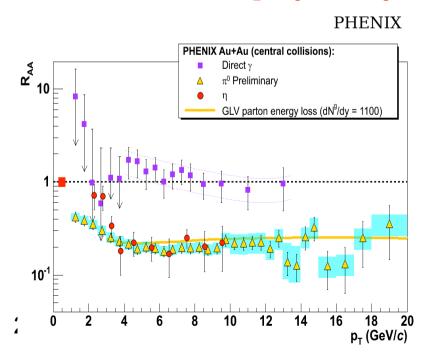


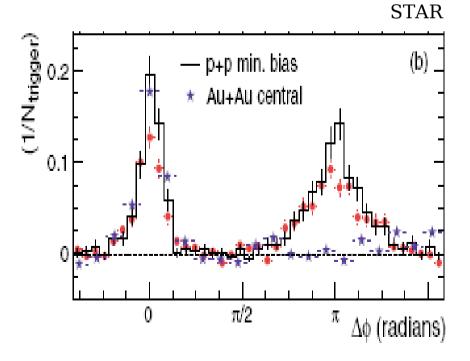


Surprises from jets in lead ion collisions

Jet quenching at RHIC

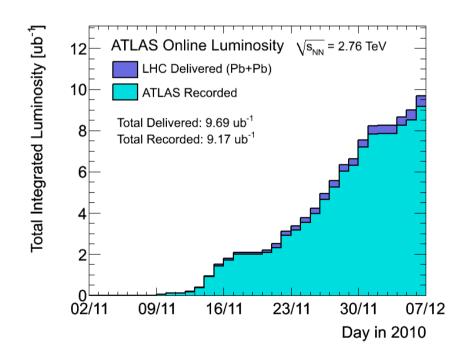
- Jet quenching was first observed in gold ion collisions at $sqrt(s_{NN}) = 100$ GeV at the Relativistic Heavy Ion Collider
 - Single hadron suppression, no photon suppression (PHENIX)
 - Dijet "disappearance" in di-hadron correlations (STAR)
 - LBL theorists & experimentalists in nuclear physics played key roles in theory, hardware, and analysis





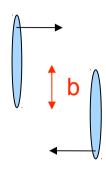
Lead ion collisions at the LHC

- LHC collided lead ions at $sqrt(s_{NN}) = 2.76$ TeV in November
- Approximately 9 ub⁻¹ recorded
 - Results shown here use 1.7 ub⁻¹, with analysis of rest of data on-going

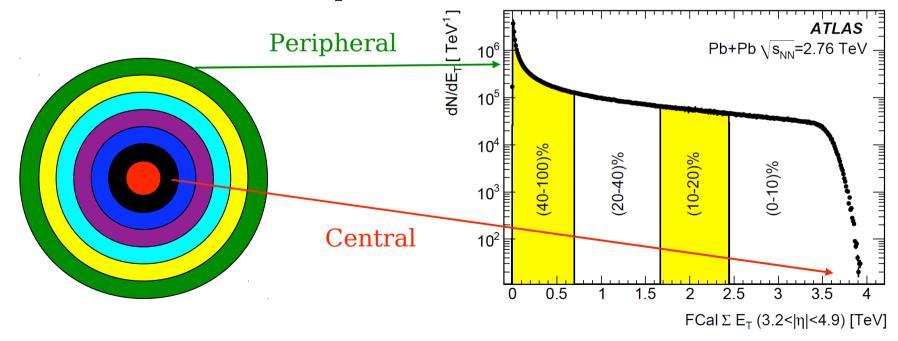


Centrality definition

- Particle multiplicity increases as classical impact parameter b decreases
- Characterize centrality by percentiles of total cross-section using Σ $E_{\scriptscriptstyle T}$ in Forward Calorimeter (FCal) spanning $3.2 < |\eta| < 4.9$



 Recover pp behavior in peripheral collisions where nuclear "overlap" is small

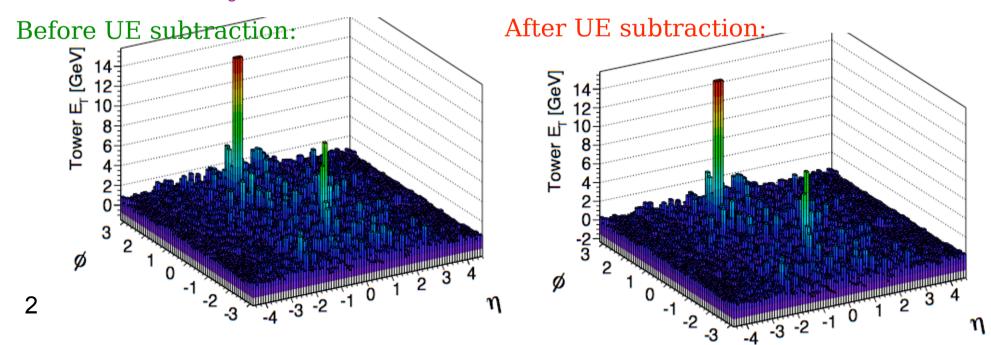


Event selection

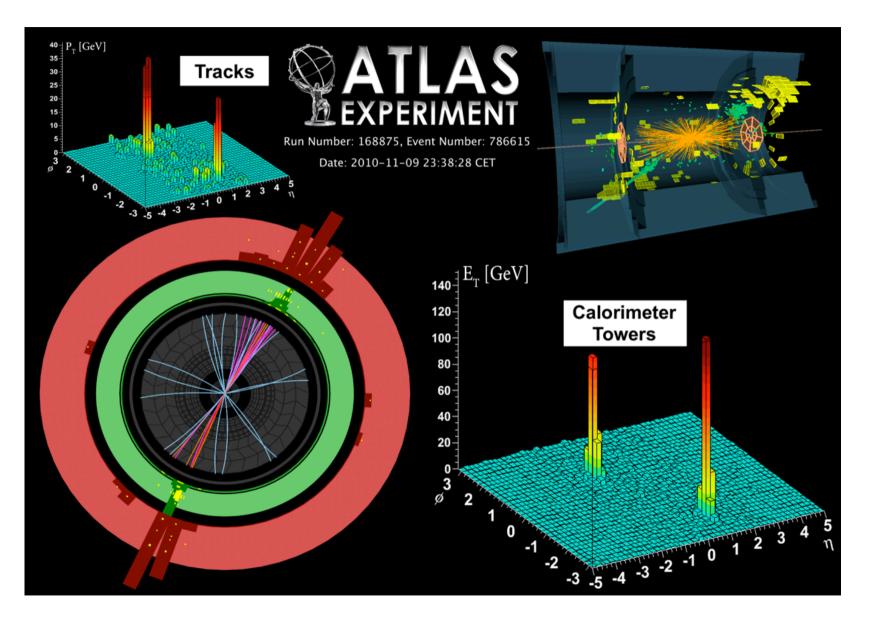
- Events triggered using minimum bias trigger
- Jets reconstructed from calorimeter towers using anti- $k_{\scriptscriptstyle T}$ algorithm with R=0.4
 - · Calibrated with energy density cell weighting
 - UE subtraction performed afterwards (next slide)
- Two leading jets are required to be within $|\eta_{1,2}| < 2.8$ where relative JES uncertainty is within 5% from pp data
- Jet $E_{T,1} > 100$ GeV so jet reconstruction in lead ion collisions is fully efficient
- Jet $E_{T,2} > 25$ GeV to be above UE background
 - · Opposite hemisphere ($\Delta \phi > \pi/2$) to select fairly "backto-back" dijet events

UE subtraction from jet

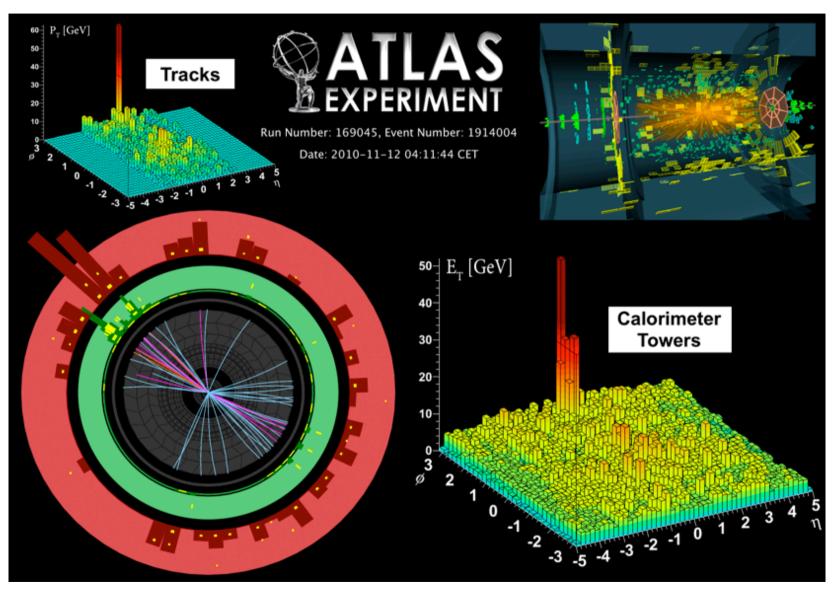
- · Underlying event in heavy ion collisions is huge \rightarrow critical to correct jet E_{T}
- · UE density is computed for each longitudinal layer, in slices of $\Delta \eta = 0.1$
 - · Average over φ (elliptical flow)
 - Exclude towers from calculation of UE density if discriminant $D = E_t^{\text{tower, max}} / E_t^{\text{tower, mean}} > 5$ (no jets are removed from analysis)
- · Jet correction is the UE density integrated against jet area
 - R=0.2 and R=0.4 used for cross-checks because UE pedestal for each jet scales with area i.e. $R^2 \rightarrow \text{results}$ are consistent



Peripheral, symmetric dijet event



Central, asymmetric dijet event

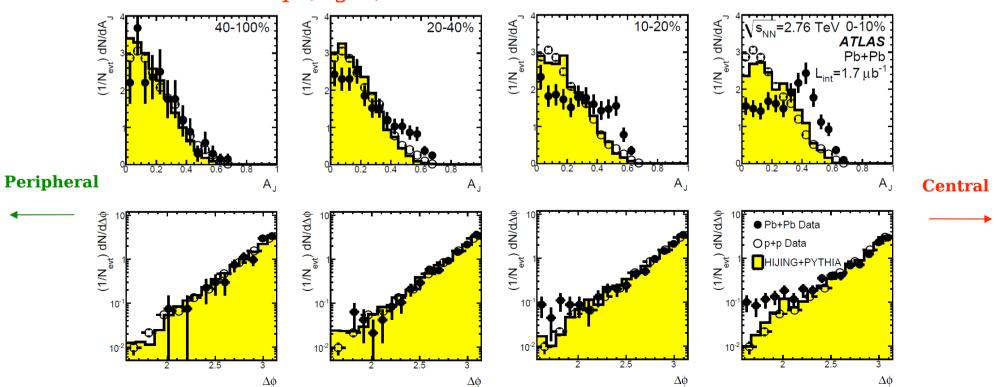


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Jet quenching at the LHC

- Asymmetry $A_{J} = (E_{T,1} E_{T,2}) / (E_{T,1} + E_{T,2})$
 - Compare Pb+Pb data to pp data and HIJING Monte Carlo (with PYTHIA dijets overlaid)
- Dijet asymmetry observed that increases with centrality
 - Most symmetric in peripheral collisions, where pp behavior is recovered (left)
 - But very asymmetric in central collisions with largest nuclear overlap (right)



Conclusions

- With $\sim\!45~\rm pb^{-1}$ of data, jet $\rm p_T\sim1.3~\rm TeV$ and dijet mass $\sim3.7~\rm TeV$ have already been observed and surpass the Tevatron reach
- First cross-section measurements of inclusive jet and dijet production at sqrt(s) = 7 TeV using 17 nb⁻¹ (accepted by Eur. Phys. J. C)
 - First measurements at hadron collider using anti- k_{T} algorithm
 - Monte Carlo based calibration scheme was developed to calibrate jets as a function of $\mathbf{p}_{\!\scriptscriptstyle T}$ and \mathbf{y}
 - JES uncertainty of 7% for central jets above 60 GeV
- New limit on dijet resonance mass: 0.87 TeV -> 1.26 TeV (published in PRL)
- New limit on contact interaction scale: 3.1 TeV -> 3.4 TeV (accepted by Phys. Lett. B)
 - QCD probed in new kinematic regime of high jet pT and large dijet mass
- First observation of jet quenching at sqrt(s_NN) = 2.76 TeV (published in PRL)

Outlook

- ATLAS jet measurements and searches will have exciting new results for the Winter 2011 conferences
 - The full 2010 dataset of 45 pb $^{\text{-}1}$ is being analyzed and will further extend the kinematic reach in both jet $p_{\text{-}1}$ and rapidity, and in dijet mass
 - Significantly reduced JES uncertainty using in-situ calibration methods is expected to be finished soon
 - Luminosity uncertainty will also be reduced
 - Other QCD analyses on dijets, multi-jets, jet shapes, etc have produced conference notes and are advancing towards publication
- Next year promises to be even more exciting with 1 fb⁻¹ or more of data
 - Steve Myers has optimistically suggested that 2-5 fb⁻¹ may be possible, so the future for jet measurements at the LHC is very bright!

Acknowledgments

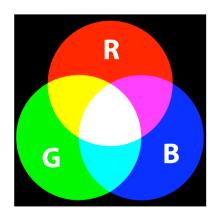


ADDITIONAL MATERIAL

Quantum chromodynamics (QCD)

- Non-abelian gauge theory where:
 - Quarks have "color", anti-quarks have "anti-color"
 - Gluons have both color and anti-color
 - Hadrons (bound states of quarks) are color neutral

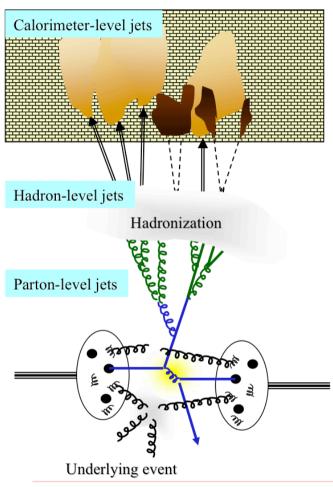




Complications with QCD

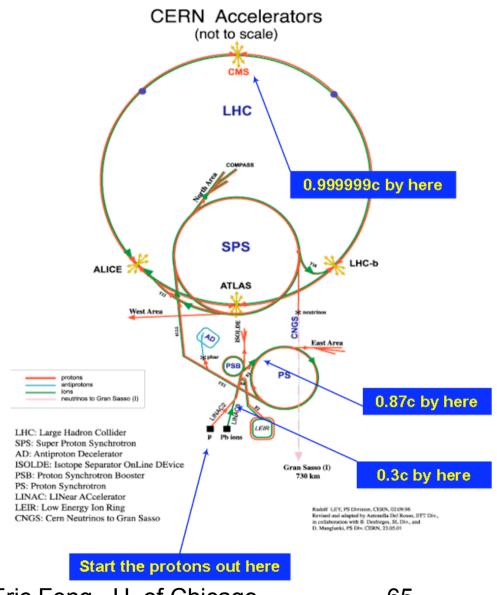
- So what's the difficulty?
- Asymptotic freedom → Partons behave as if free at large Q² i.e. very short distances
- But QCD becomes non-perturbative at low Q² (long distances)
 - Quark confinement
 - We measure jets (collimated flows of hadrons), not partons





 This talk discusses how we have performed jet measurements in the ATLAS experiment and used these to probe *partonic* predictions of NLO pQCD, as well as to search for **new physics**

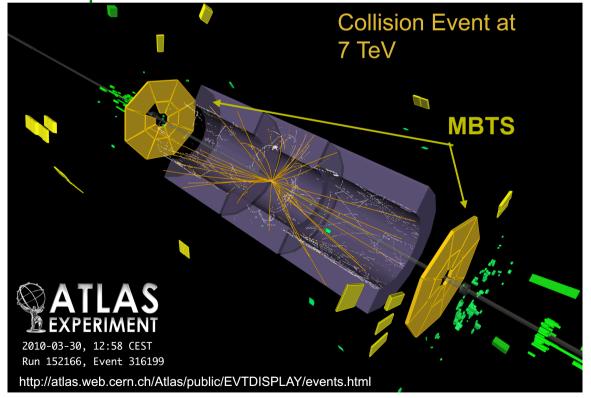
LHC accelerator complex



MBTS Trigger

- MBTS (Minimum Bias Trigger Scintillator) inclusive trigger
 - Require at least one scintillator fired from either η hemisphere: $2.09 < |\eta| < 3.84$

No significant bias introduced to the inclusive jet sample



Online and offline selection

- Only summarize briefly below due to time constraints MUCH more detail in backup slides!
- Data sample studied is between 17 300 nb⁻¹ of data at 7 TeV collected through Period D before ICHEP
- Official Jet/EtMiss Good Run Lists and require "green" flag for luminosity
- Require at least one vertex reconstructed within |z| < 10cm of detector center to suppress beam halo and beam gas
- Require L1_J5 trigger fired from L1Calo stream
 - Restrict leading (sub-leading) jet to p_T>60 GeV (30 GeV) so that jet trigger and reconstruction are both fully efficient

What is the anti-k_T jet algorithm?

- Cone jet algorithms (sum all energy in some region)
 - ATLAS Cone (seeded → not infrared nor collinear safe)
 - SIS-Cone (Seedless Infrared Safe)
- Clustering AKA "sequential recombination" jet algorithms (invert the radiation and fragmentation)

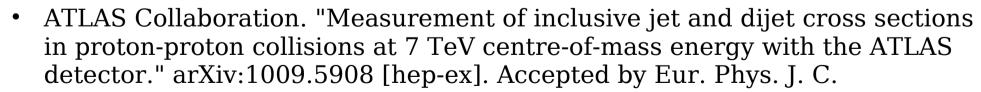
$$d_{ij} = \min(k_{ti}^{2p}, k_{tj}^{2p}) \frac{\Delta_{ij}^2}{R^2} \qquad \Delta_{ij}^2 = (y_i - y_j)^2 + (\phi_i - \phi_j)^2$$

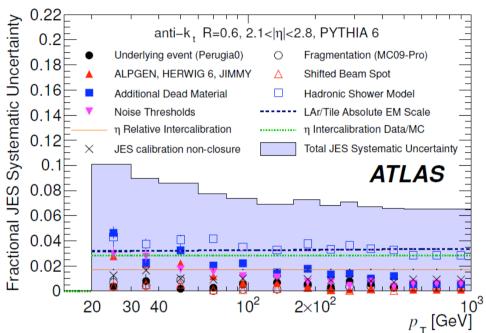
- K_{T} (p= 1): Clusters softest constituents first
- Cambridge (p= 0): Clusters closest constituents first
- Anti-k_⊤ (p=-1): Clusters hardest constituents first
- → Reconstructs the same jets no matter whether there's soft radiation near the initial constituents

Absolute JES uncertainty

• Absolute JES uncertainty is slightly larger for forward jets and significantly bigger at lower jet $p_{\scriptscriptstyle T}$

- Up to 10% at $p_{\scriptscriptstyle T}$ = 20 GeV
- Dominant sources:
 - LAr/Tile EM scale: 3%
 - Hadronic shower: ~2-4%
 - Cluster noise thresholds
 - Material description
 - η intercalibration
 - Physics: Alpgen vs. Pythia



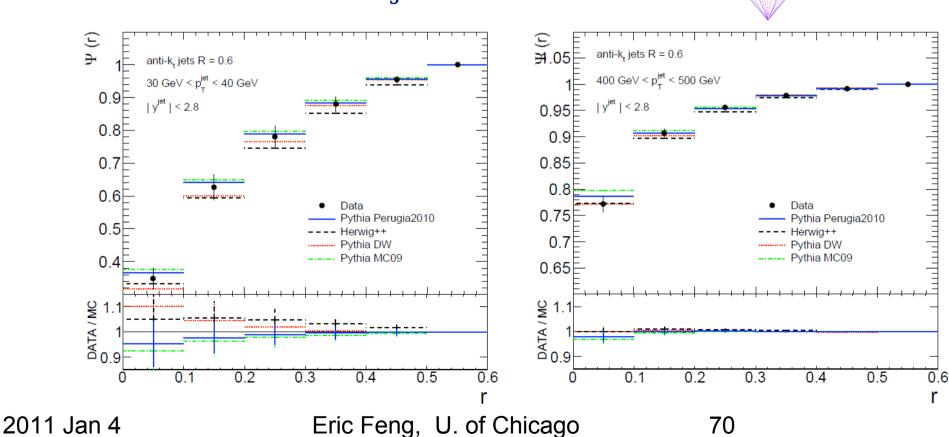


Integral jet shape method

 $\Psi(r)$

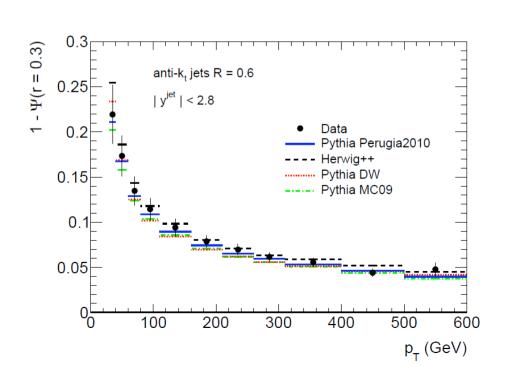
R

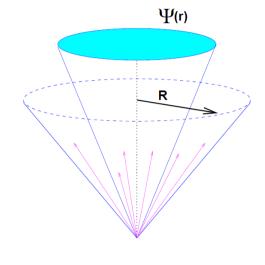
- Integral jet shapes can be computed as integral of the differential jet shape from the jet axis out to some radius R
 - Dominated by transverse momentum near jet axis



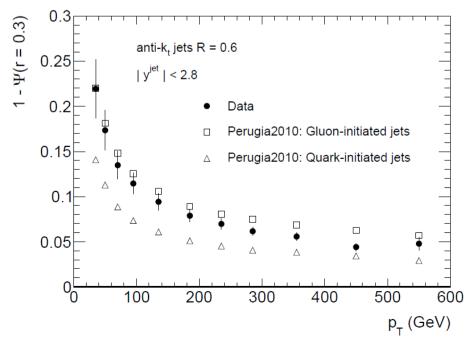
Integral jet shape

- Integral jet shapes $\Psi(r)$ can be computed as integral of the differential jet shape from the jet axis out to some radius R
 - Dominated by measurement nearest to jet axis
- Thus 1 Ψ (r=0.3) is the fraction of transverse momentum outside of r=0.3
- May eventually be useful to separate gluon and quark jets









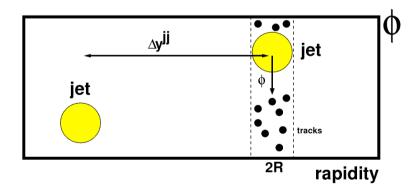
Charged particle flow: Method

Charged particle flow in inclusive dijet events:

$$<\frac{d^2p_T}{|d\phi|dy}>_{jets} = \frac{1}{2R|\Delta\phi|} \frac{1}{N^{jet}} \sum_{jets} p_T(|\phi - \Delta\phi/2|, |\phi + \Delta\phi/2|), \text{ with } 0 \le |\phi| \le \pi$$

is the average transverse momentum as a function of the azimuthal distance from the jet axis and rapidity separation between the two leading jets

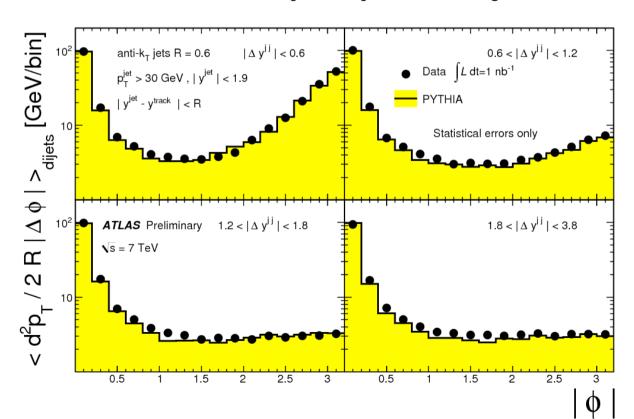
• Here p_T is computed as the scalar sum of the transverse momenta of tracks at a given angle ϕ with respect to jet axis



- Only tracks within rapidity range occupied by jet are used
- Jet required to be within |y|<1.9 to ensure that jet is fully within tracker acceptance |y|<2.5
- Track-based method is useful to confirm results from calorimeter-based jet shapes

Charged particle flow: Data distributions

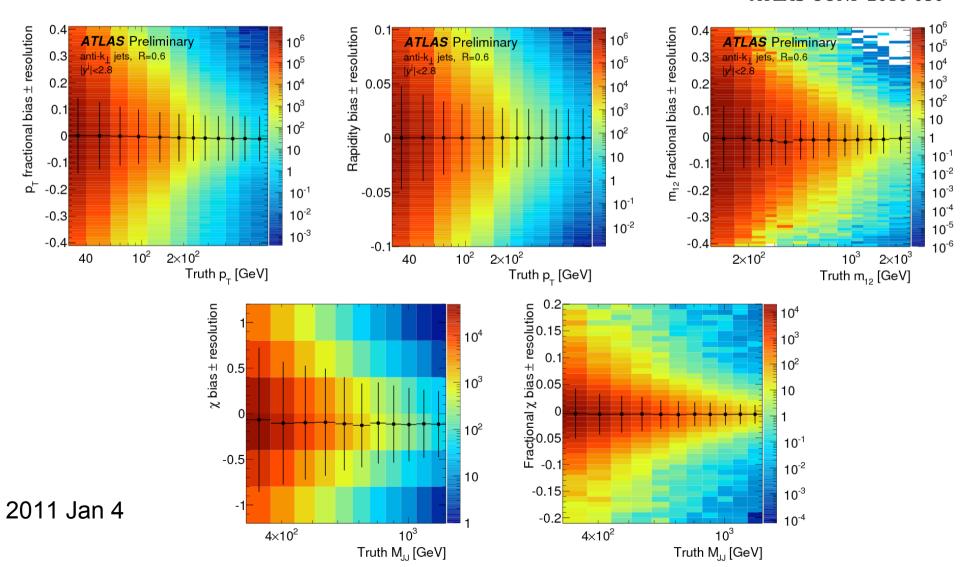
- For $|\Delta y^{ij}| < 0.6$, two collimated flows of charged particles (dijets) observed at $|\phi|=0$ and $|\phi|=\pi$
- For $|\Delta y^{ij}| > 1.2$, jet structure observed at low $|\phi|$ followed by plateau of remaining hadronic activity as $|\phi|$ increases
- Monte Carlo provides reasonable description of data, but slightly underestimates hadronic activity away from the jet direction (see backup)



Resolutions of jet observables

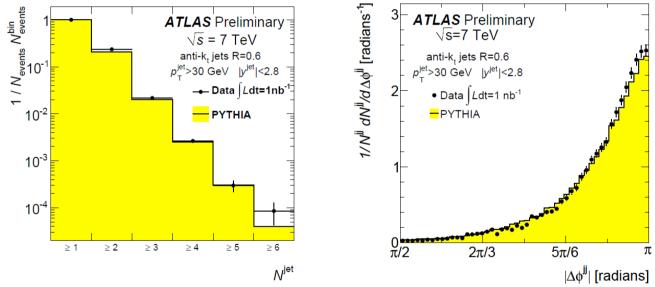
• Jet energy and angular resolutions studied in Monte Carlo in order to assign appropriate bin widths

ATLAS-CONF-2010-050



Jet observation at sqrt(s) = 7 TeV

- With 1 nb⁻¹, reported observation of jets at center-of-mass energy of 7 TeV on behalf of Jet/EtMiss and Standard Model groups
- Measured inclusive jet pT spectrum, dijet mass spectrum, and dijet azimuthal decorrelation

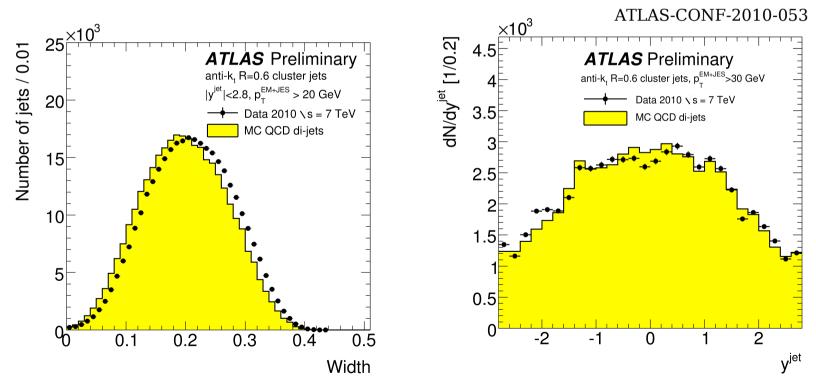


• E. Feng (for the ATLAS Collaboration). "Observation of Energetic Jet Production in pp Collisions at sqrt(s) = 7 TeV using the ATLAS Experiment at the LHC." arXiv:1010.1974 [hep-ex]. To appear in *Proceedings of PLHC 2010*, Hamburg, Germany, June 2010.

Jet properties

Jet width and rapidity

- Studied inclusive jet, dijet, and multi-jet observables to test description of calorimeters, jets, and missing $E_{\scriptscriptstyle T}$ in Monte Carlo
- Jet radial width shown to be wider in data than predicted by MC09
 - Consistent with jet shapes!
- Monte Carlo description was shown to be satisfactory for jets in |y|<2.8



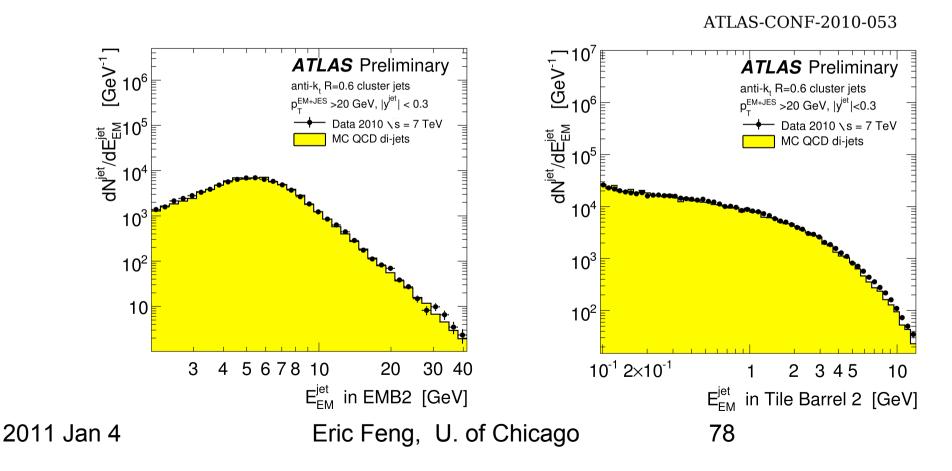
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Jet energy in different calorimeter layers

- Detailed studies of jet energy deposited in each layer of calorimeter
- Second layer of EM and hadronic calorimeters shown below
 - Good agreement between data and MC

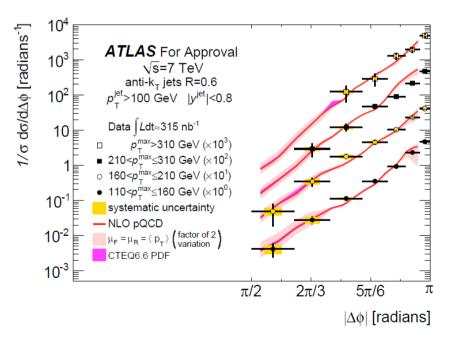


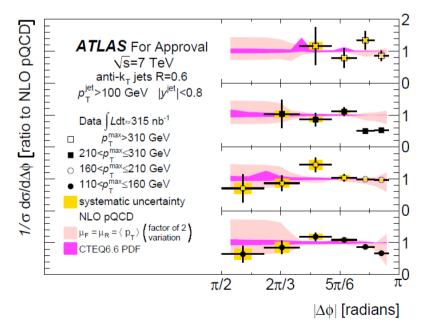
More tests of perturbative QCD

Dijet Azimuthal Decorrelation ($\Delta \phi$)

- Dijet angular distribution: $\Delta \phi = |\phi_1 \phi_2|$ measured in bins of leading jet p_T
- Peak at $\Delta \phi = \pi \rightarrow$ dominant final state is back-to-back dijets
 - Deviation from $\Delta \phi = \pi$ due to radiation of one or more gluons
 - 3-jet final state calculated using NLO pQCD

ATLAS-COM-CONF-2010-080



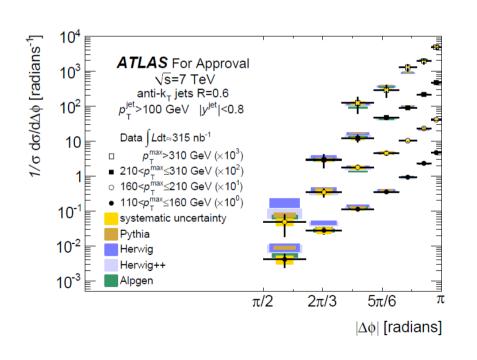


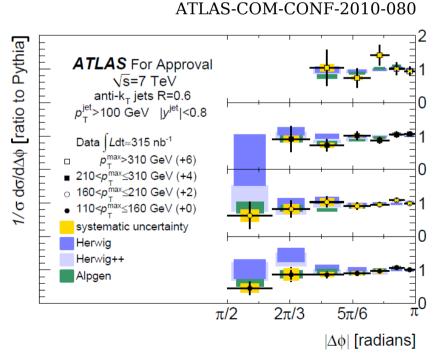
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Dijet Azimuthal Decorrelation ($\Delta \phi$)

- Four or more real emissions is best modeled using a leading-order ME generator with 2->N matrix elements
 - Sensitive to different angular distributions produced by 2->2 vs. 2->N parton shower Monte Carlos
- Useful to tune amount of ISR/FSR and underlying event in parton shower Monte Carlo

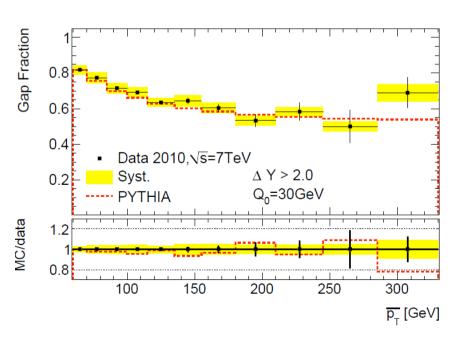


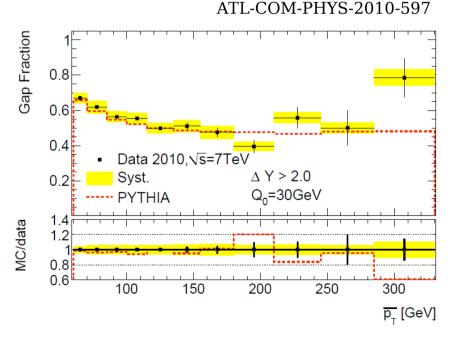


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Dijets with Rapidity Gaps

- Measure fraction of events with gap between two boundary jets:
 - · Two highest p_T jets (left), OR
 - Most forward and most backward jet (right)
- Sensitive to BFKL dynamics vs. DGLAP evolution
- · Also can study wide-angle soft-gluon radiation
- Uses Monte-Carlo based JES uncertainty up to |y|<4.5



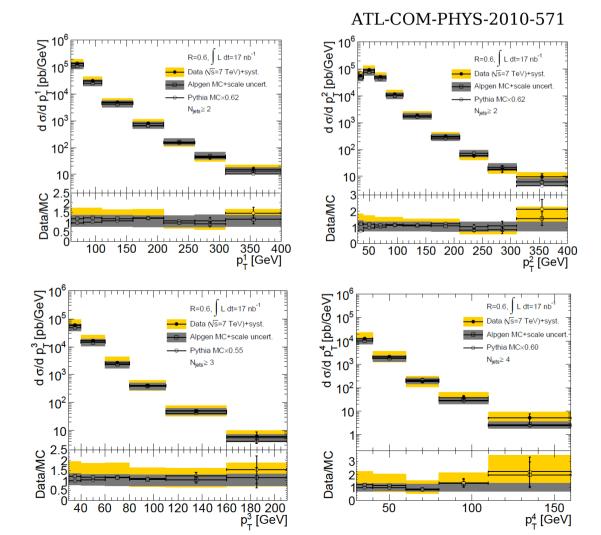


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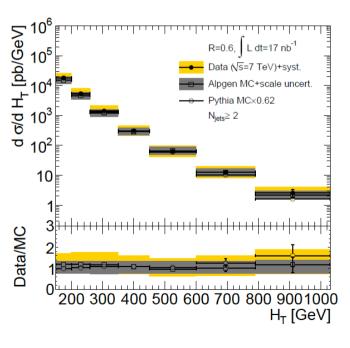
$Multi-Jets: \\ p_{\scriptscriptstyle T}\text{-ordered cross-sections}$

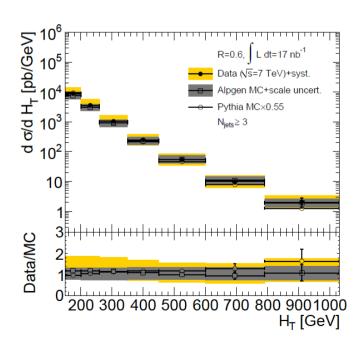
- Subdivision of the inclusion jet cross-section into separate crosssections for each of the $p_{\scriptscriptstyle T}$ -ordered jets
- Sensitive to matrix element and final state radiation

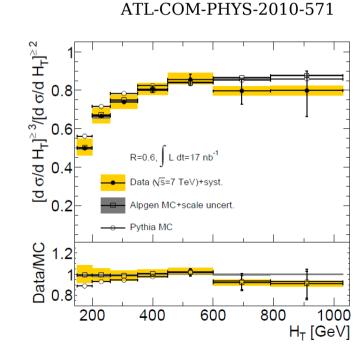


Multi-Jets – H_T

- \cdot H_T = scalar sum of jet p_T
- Ratio of 3-jet and 2-jet cross-sections in H_T is a direct probe of α_S
- Additional uncertainty in absolute JES due to flavor (quark vs. gluon) as well as close-by jets has been studied







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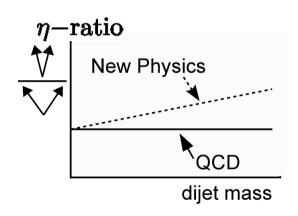
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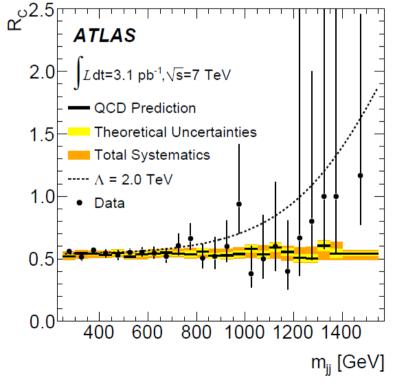
Even more QCD analyses

- Many more jet and QCD-related analyses finished or in progress
 - Inclusive jet cross-section measured using trackjets
 - Jet fragmentation measured using tracks
 - Vector boson production in association with jets (W/Z + jets)
 - Diffractive dijets
 - Etc...
- All of these analyses are very interesting, but unfortunately do not have time to describe them here...

Search for contact interactions using the η ratio

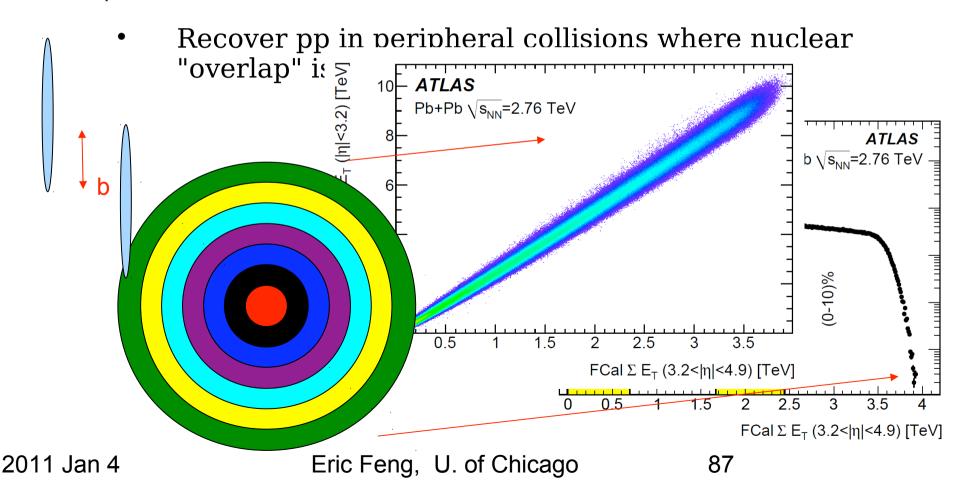
- Checked ratio of large-opening to small-opening jets for tail as a function of dijet mass: η -ratio = $\frac{N(|\eta_{1,2}|<0.5)}{N(0.5<|\eta_{1,2}|<1)}$
 - Like χ , largely cancels out JES uncertainty and leaves relative JES
 - Also sensitive to contact interactions from small-angle scattering



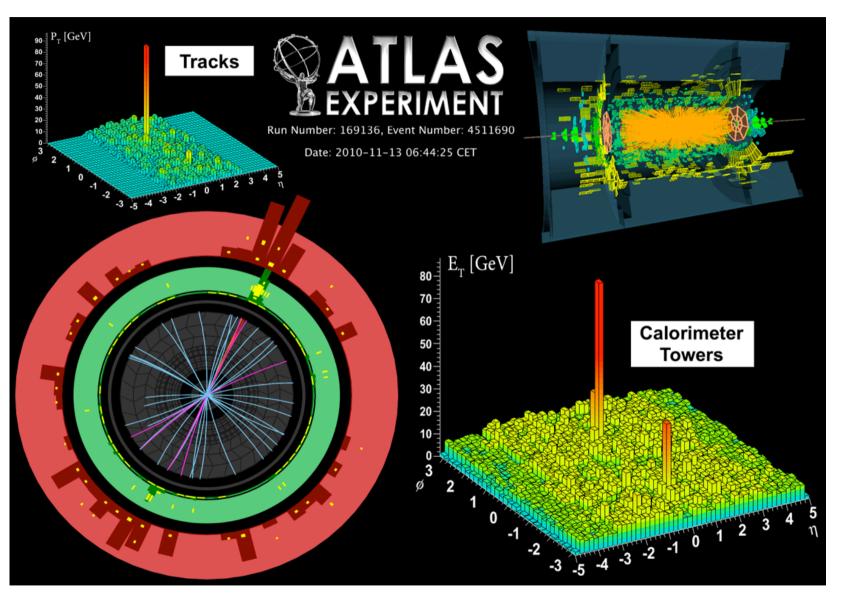


Centrality definition

- Particle multiplicity increases as classical impact parameter b decreases
- Characterize centrality by percentiles of total cross-section using $\Sigma E_{\scriptscriptstyle T}$ in Forward Calorimeter (Fcal) spanning 3.2 < $|\eta| < 4.9$



More central, asymmetric dijet event



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Central event, with split dijet and additional activity

